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Some Differentials in the mod 3 Adams Spectral Sequence

by

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Let A_p be the mod p Steenrod algebra. J. F. Adams [1] introduced a spectral sequence which has as its E_2 term $Ext_{A_p}(H*(X), Z_p)$ and which converges to a graded algebra associated to $\pi_{\bullet}^{s}(X;p)$, i. e., the p-primary stable homotopy groups of X. In this paper we will study this sequence for $X=S^n$, p=3. The first problem in any use of the Adams spectral sequence is to obtain $E_2 = Ext_{A_3}^{s,t}(Z_3, Z_3)$. We do this by the technique of J. P. May [5]. J. P. May constructed another spectral sequence which has as its E_{∞} term an algebra $E^{\circ}Ext$, i. e. a tri-graded algebra associated to $E_2 = Ext$. In [9], we extended (and corrected) May's computations to obtain complete information on EoExt through dimension 158. The next problem is to obtain the differentials in the Adams spectral sequence. J. P. May [5] and S.Oka have previously determined all differentials at least in the range t-s = 77 by using the results of the 3 - components of stable homotopy groups of sphere which have been calculated by H.Toda [12,13,14,15], J.P.May [5] and S.Oka [10]. The purpose of this paper is to evaluate the differentials in the range 78 ≤t-s≤104. Our main result is Theorem 3.19.

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\S 1 Algebra Structure of $Ext^{**}_{A_3}(Z_3, Z_3)$

From now on we will write $H^{**}(A_3)$ instead of $Ext^{**}_{A_3}$ (Z_3, Z_3) for the E_2 term of the Adams spectral sequence. The table of $H^{**}(A_3)$ which will be needed for this paper is given in Appendix. Relations involving a_0 and h_0 are indicated by vertical and slanting lines respectively. Since we have computed $H^{**}(A_3)$ by May's techniques, the products which we naturally obtained are actually the products according to the algebra structure of $E^0H^{**}(A_3)$. The product in $H^{**}(A_3)$ of

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two elements always contains as a summand their product in $E \circ H^{**}(A_3)$ but may possibly contain also other terms of the same bi-grading (s,t) but of lower weight in the sense of J. P. May [5]. The relations holding in $E \circ H^{**}(A_3)$ which cannot be listed for reasons of space are easily obtained from [9. Theorem 3.3 and 4.3].

The following relations in $H^{**}(A_3)$ differ from ones in $E^0H^{**}(A_3)$. This list is by no means complete.

PROPOSITION 1.1. We have the following relations in $H^{**}(A_3)$.

a).
$$u h_2 = -a_0^2 c$$
 (58 stem)

b).
$$h_0 \cdot h_1 b_{02} = -b_{11}k$$
 (60 stem)

c).
$$h_1 \cdot h_1 b_{02} = b_{11}^2$$
 (68 stem)

d).
$$h_1 h_2 b_{02} = b_{01} d - c b_{11}$$
 (92 stem)
 $h_2 h_1 b_{02} = b_{01} d - c b_{11}$

e).
$$b_{02}h_2 g_2 = b_{01}v_0$$
 (99 stem)

f).
$$h_2 h_0 u b_{02} = -h_0 b_{01} G$$
 (107 stem)

PROOF. We first consider the relation b). Let $\overline{b_{02}}$ be a cochain

$$\begin{array}{l} [\,\varepsilon_{\,2}\,|\,\,\varepsilon_{\,2}^{\,2}\,] \,+\, [\,\varepsilon_{\,2}^{\,2}\,|\,\,\varepsilon_{\,2}\,] \,+\, [\,\varepsilon_{\,1}^{\,3}\,|\,\,\varepsilon_{\,1}\,\varepsilon_{\,2}^{\,2}\,] \,+\, [\,\varepsilon_{\,1}^{\,3}\,\varepsilon_{\,2}^{\,2}\,|\,\,\varepsilon_{\,1}\,] \\ -\, [\,\varepsilon_{\,1}^{\,3}\,\varepsilon_{\,2}\,|\,\,\varepsilon_{\,1}\,\varepsilon_{\,2}\,] \,+\, [\,\varepsilon_{\,1}^{\,6}\,\varepsilon_{\,2}\,|\,\,\varepsilon_{\,2}^{\,2}\,] \,+\, [\,\varepsilon_{\,1}^{\,6}\,\varepsilon_{\,2}\,|\,\,\varepsilon_{\,2}^{\,2}\,] \end{array}$$

in the cobar construction $F^*(A_3^*)$. Let $\overline{b_{01}}$ and $\overline{b_{11}}$ be cocycles

$$[\xi_1 \mid \xi_1^2] + [\xi_1^2 \mid \xi_1]$$
 and $[\xi_1^3 \mid \xi_1^6] + [\xi_1^6 \mid \xi_1^3]$

in $F^*(A_3^*)$ which represent the elements $b_{01} \in H^{2,12}$ (A_3) and $b_{11} \in H^{2,36}(A_3)$, respectively. By routine calculations, we see that $b_{02}h_1 \in H^{3,60}(A_3)$ is represented by a cocycle $\overline{b_{02}h_1}$ in $F^*(A_3^*)$:

$$\begin{array}{l} \overline{b_{02}h_1} = \overline{b_{02}} \cdot \left[\begin{smallmatrix} \xi & 3 \\ 1 \end{smallmatrix} \right] - \overline{b_{11}} \cdot \left[\begin{smallmatrix} \xi & 6 \\ 1 \end{smallmatrix} \right] - \left[\begin{smallmatrix} \xi & 9 \\ 1 \end{smallmatrix} \right] \left[\begin{smallmatrix} \xi & 1 \\ 1 \end{smallmatrix} \right] \\ - \left[\begin{smallmatrix} \xi & 9 \\ 1 \end{smallmatrix} \right] \left[\begin{smallmatrix} \xi & 2 \\ 1 \end{smallmatrix} \right] \left[\begin{smallmatrix} \xi & 4 \\ 1 \end{smallmatrix} \right] - \left[\begin{smallmatrix} \xi & 9 \\ 1 \end{smallmatrix} \right] \left[\begin{smallmatrix} \xi & 1 \\ 1 \end{smallmatrix} \right] - \left[\begin{smallmatrix} \xi & 9 \\ 1 \end{smallmatrix} \right] \left[\begin{smallmatrix} \xi & 5 \\ 1 \end{smallmatrix} \right] \left[\begin{smallmatrix} \xi & 5 \\ 1 \end{smallmatrix} \right] - \left[\begin{smallmatrix} \xi & 3 \\ 2 \end{smallmatrix} \right] \cdot \overline{b_{01}} \cdot \end{array}$$

By routine calculations, we have that

$$\begin{split} \delta & \{ -\overline{b_{02}} \cdot [\varepsilon_2] + [\varepsilon_1^9 \mid \varepsilon_1^2 \mid \varepsilon_1 \varepsilon_2] + [\varepsilon_1^9 \mid \varepsilon_1 \mid \varepsilon_1^2 \varepsilon_2] + [\varepsilon_1^9 \mid \varepsilon_1 \varepsilon_2 \mid \varepsilon_1^2] \\ & + [\varepsilon_1^9 \mid \varepsilon_1^2 \varepsilon_2 \mid \varepsilon_1] + [\varepsilon_1^9 \mid \varepsilon_1^5 \mid \varepsilon_1^2] - [\varepsilon_2^3 \mid \varepsilon_1^2 \mid \varepsilon_1^2] + [\varepsilon_3 \mid \varepsilon_1^2 \mid \varepsilon_1] \\ & + [\varepsilon_3 \mid \varepsilon_1 \mid \varepsilon_1^2] + [\varepsilon_1^9 \mid \varepsilon_1^4 \mid \varepsilon_1^3] + [\varepsilon_1^9 \mid \varepsilon_1^3 \mid \varepsilon_1^4] \} \\ & = \overline{b_{02} h_1} \cdot [\varepsilon_1] - \overline{b_{11}} \cdot ([\varepsilon_1^3 \mid \varepsilon_2] - [\varepsilon_1^6 \mid \varepsilon_1]). \end{split}$$

Since it is easy to see that the cocycle $\begin{bmatrix} \xi_1^3 \mid \xi_2 \end{bmatrix} - \begin{bmatrix} \xi_1^6 \mid \xi_1 \end{bmatrix}$ represents the element $k \in H^{2,28}(A_2)$, we have the relation b).

Next we consider the relation c). By tedious but routine calculations, we have

$$\begin{array}{lll} \delta & \{ \; \overline{b_{02}} \cdot [\; \xi_{\; 1}^6] \; - \; [\; \xi_{\; 1}^9 \; | \; \xi_{\; 1}^2 \; | \; \xi_{\; 1}^7] \; - \; [\; \xi_{\; 1}^9 \; | \; \xi_{\; 1}^8] \; + \; [\; \xi_{\; 1}^9 \; | \; \xi_{\; 1}^5 \; | \; \xi_{\; 1}^4] \\ & + \; [\; \xi_{\; 1}^9 \; | \; \xi_{\; 1}^4 \; | \; \xi_{\; 1}^5] \; - [\; \xi_{\; 1}^9 \; | \; \xi_{\; 1}^8 \; | \; \xi_{\; 1}^7] \; - [\; \xi_{\; 1}^9 \; | \; \xi_{\; 1}^7] \; + [\; \xi_{\; 2}^3 \; | \; \xi_{\; 1}^4 \; | \; \xi_{\; 1}^4] \\ & + \; [\; \xi_{\; 2}^3 \; | \; \xi_{\; 1}^5] \; + [\; \xi_{\; 2}^3 \; | \; \xi_{\; 1}^5 \; | \; \xi_{\; 1}^7] \; + [\; \xi_{\; 2}^3 \; | \; \xi_{\; 1}^4 \; | \; \xi_{\; 2}^7] \\ & + \; [\; \xi_{\; 3}^3 \; | \; \xi_{\; 1}^3 \; | \; \xi_{\; 1}^2] \; - [\; \xi_{\; 2}^3 \; | \; \xi_{\; 1}^8 \; | \; \xi_{\; 1}^7] \; + [\; \xi_{\; 2}^3 \; | \; \xi_{\; 1}^4 \; | \; \xi_{\; 1}^7] \\ & + \; [\; \xi_{\; 1}^9 \; | \; \xi_{\; 1}^2] \; + [\; \xi_{\; 1}^9 \; | \; \xi_{\; 1}^3 \; \xi_{\; 2} \; | \; \xi_{\; 1}^7] \} \\ & = \; \overline{b_{02}h_1} \; \cdot \; [\; \xi_{\; 3}^3 \; | \; + \; \overline{b_{11}} \cdot \overline{b_{11}} \; \cdot \end{array}$$

Then we have the relation c).

Next we consider the relation d). By routine calculations, we see that h_2b_{02} ϵ $H^{3,84}$ (A_3) is represented by a cocycle $\overline{h_2 b_{02}}$ in $F^*(A_3^*)$:

$$\overline{h_2\,b_{02}} = \left[\,\xi\, {}_1^9 \right] \cdot \overline{b_{02}} \, - \, \left[\,\xi\, {}_1^{18} \right] \cdot \overline{b_{01}} \, - \, \left[\,\xi\, {}_2^3 \right] \cdot \overline{b_{11}} \, + \, \left[\,\xi\, {}_1^9 \, \right] \,\xi\, {}_1^6 \, \right].$$
 By routine calculations, we have

$$\begin{array}{lll} \delta & (\ -\ [\ \xi \ _1^{12}\] \ \cdot \ b_{02} \ +\ [\ \xi \ _2^{3}\] \ \cdot \ \overline{b_{02}} \ +\ [\ \xi \ _1^{9}\ \xi \ _2^{3}\] \ \cdot \ \overline{b_{01}} \ +\ [\ \xi \ _1^{21}\] \ \cdot \ \overline{b_{01}} \\ & +\ [\ \xi \ _1^{3}\ \xi \ _2^{3}\] \ \cdot \ \overline{b_{11}} \ -\ [\ \xi \ _1^{12}\] \ \xi \ _1^{6}\] \ +\ [\ \xi \ _1^{3}\] \) \ \cdot \ \overline{b_{01}} \\ & =\ [\ \xi \ _1^{3}\] \ \cdot \ \overline{b_{2}} \ -\ (\ [\ \xi \ _1^{9}\] \ \xi \ _2^{3}\] \ -\ [\ \xi \ _1^{18}\] \ \xi \ _1^{3}\] \) \ \cdot \ \overline{b_{01}} \\ & +\ (\ [\ \xi \ _2^{3}\] \ \xi \ _1^{3}\] \ -\ [\ \xi \ _1^{9}\] \ \xi \ _1^{6}\] \) \ \cdot \ \overline{b_{11}} \ . \end{array}$$

Since it is easy to see that the cocycles $\begin{bmatrix} \xi \ _1^9 \mid \xi \ _2^3 \end{bmatrix} - \begin{bmatrix} \xi \ _1^{18} \mid \xi \ _1^{3} \end{bmatrix}$ and $\begin{bmatrix} \xi \ _2^3 \mid \xi \ _1^3 \end{bmatrix} - \begin{bmatrix} \xi \ _1^9 \mid \xi \ _1^6 \end{bmatrix}$ represent the elements $d \in H^{2,84}(A_3)$ and $c \in H^{2,60}(A_3)$, respectively, we have the relation $h_1 \cdot h_2 b_{02} = b_{01} d - c b_{11}$. Similarly, we have

$$\begin{split} \delta & (- \overline{b_{02}} \cdot [\, \xi_{\,1}^{\,12}] \, + \, \overline{b_{02}} \cdot [\, \xi_{\,2}^{\,2}] \, + \, \overline{b_{11}} \cdot [\, \xi_{\,1}^{\,15}] \, + \, \overline{b_{11}} \cdot [\, \xi_{\,1}^{\,3} \xi_{\,2}^{\,2}] \\ & - [\, \xi_{\,1}^{\,9}] \, |\, \xi_{\,1}^{\,2} \, |\, \xi_{\,1}^{\,2} \xi_{\,2}^{\,2}] \, - [\, \xi_{\,1}^{\,9}] \, |\, \xi_{\,1}^{\,2} \xi_{\,2}^{\,2}] \, - [\, \xi_{\,1}^{\,9}] \, |\, \xi_{\,1}^{\,2} \xi_{\,2}^{\,2}] \, |\, \xi_{\,1}^{\,1}] \\ & - [\, \xi_{\,1}^{\,9}] \, |\, \xi_{\,1}^{\,2} \xi_{\,2}^{\,2}] \, |\, \xi_{\,1}^{\,2}] \, + [\, \xi_{\,2}^{\,3}] \, |\, \xi_{\,1}^{\,2}] \, + [\, \xi_{\,2}^{\,3}] \, |\, \xi_{\,1}^{\,1}] \, + [\, \xi_{\,2}^{\,3}] \, |\, \xi_{\,1}^{\,1}] \, + [\, \xi_{\,1}^{\,9}] \, |\, \xi_{\,1}^{\,2}] \, + [\, \xi_{\,1}^{\,9} \xi_{\,2}^{\,3}] \, \cdot \, \overline{b_{01}} \\ & + [\, \xi_{\,1}^{\,9}] \, |\, \xi_{\,1}^{\,1}] \, |\, \xi_{\,1}^{\,14}] \, + [\, \xi_{\,1}^{\,9}] \, |\, \xi_{\,1}^{\,2}] \, + [\, \xi_{\,1}^{\,9}] \, |\, \xi_{\,1}^{\,14} \, |\, \xi_{\,1}] \\ & + [\, \xi_{\,1}^{\,9}] \, |\, \xi_{\,1}^{\,13}] \, |\, \xi_{\,1}^{\,2}] \, + [\, \xi_{\,1}^{\,9}] \, |\, \xi_{\,1}^{\,4}] \, |\, \xi_{\,1}^{\,11}] \, + [\, \xi_{\,1}^{\,9}] \, |\, \xi_{\,1}^{\,5}] \, |\, \xi_{\,1}^{\,10}] \,) \\ & = \overline{b_{02}h_{1}} \cdot [\, \xi_{\,1}^{\,9}] \, - \, \overline{b_{11}} \cdot (\, [\, \xi_{\,2}^{\,3}] \, |\, \xi_{\,1}^{\,3}] \, - \, [\, \xi_{\,1}^{\,9}] \, |\, \xi_{\,1}^{\,6}] \,) \\ & + (\, [\, \xi_{\,1}^{\,9}] \, |\, \xi_{\,2}^{\,3}] \, - \, [\, \xi_{\,1}^{\,18}] \, |\, \xi_{\,1}^{\,3}] \,) \cdot \, \overline{b_{01}} \, . \end{split}$$

Then we have the relation $b_{02}h_1 \cdot h_2 = b_{11} c - db_{01}$.

Next we consider the relation e). Let α , β and γ be cochains

$$-[\xi_{1}^{20}]\xi_{1}],$$

respectively. Then $b_{02}h_2 \in H^{3.84}(A_3)$ is represented by a cocycle $\overline{b_{02}h_2}$ in $F^*(A_3^*)$: $\overline{b_{02}h_2} = \overline{b_{02}} \cdot [\xi_1^9] + \overline{b_{11}} \cdot [\xi_1^{12}] - \overline{b_{11}} \cdot [\xi_2^3] + [\xi_1^{18}] \cdot \overline{b_{01}} - [\xi_1^9] \cdot \alpha$.

By routine calculations, we have

$$\begin{split} \delta & \{ \, \overline{b_{02}} \, \, . \, \, (\, - \, [\, \varepsilon_{\, 3} \, | \, \varepsilon_{\, 1} \,] \, + \, [\, \varepsilon_{\, 2}^{\, 3} \, | \, \varepsilon_{\, 1}^{\, 2} \,] \,) \, - \, \overline{b_{11}} \, \, . \, \, \beta \\ & - \, \alpha \, . \, (\, - \, [\, \varepsilon_{\, 3} \, | \, \varepsilon_{\, 1} \,] \, + \, [\, \varepsilon_{\, 2}^{\, 3} \, | \, \varepsilon_{\, 2}^{\, 2} \,] \,) \, + \, \, \gamma \, . \, (\, [\, \varepsilon_{\, 2} \, | \, \varepsilon_{\, 1} \,] \, - \, [\, \varepsilon_{\, 1}^{\, 3} \, | \, \varepsilon_{\, 1}^{\, 2} \,] \,) \\ & = \, \overline{b_{02}h_2} \, . \, \, (\, [\, \varepsilon_{\, 2} \, | \, \varepsilon_{\, 1} \,] \, - \, [\, \varepsilon_{\, 1}^{\, 3} \, | \, \varepsilon_{\, 1}^{\, 2} \,] \,) \\ & - \, \overline{b_{01}} \, . \, \, (\, - \, [\, \varepsilon_{\, 1}^{\, 9} \, | \, \varepsilon_{\, 3} \, | \, \varepsilon_{\, 1}^{\, 2} \,] \,) \\ & - \, [\, \varepsilon_{\, 1}^{\, 18} \, | \, \varepsilon_{\, 2}^{\, 3} \, | \, \varepsilon_{\, 1}^{\, 2} \,] \,) \, . \end{split}$$

It is easy to see that

$$[\xi_{2} | \xi_{1}] - [\xi_{1}^{3} | \xi_{1}^{2}]$$

and

 $-\left[\xi_{1}^{9}\mid\xi_{3}\mid\xi_{1}\right]+\left[\xi_{1}^{9}\mid\xi_{2}^{3}\mid\xi_{1}^{2}\right]+\left[\xi_{1}^{18}\mid\xi_{2}\mid\xi_{1}\right]-\left[\xi_{1}^{18}\mid\xi_{1}^{3}\mid\xi_{1}^{2}\right]$ are representatives of $g_{2}\in H^{2,20}(A_{3})$ and $v_{0}\in H^{3,92}(A_{3})$, respectively. Then we have the relation e).

Next we consider the relation a). It is easy to see that $u \in H^{3,26}(A_3)$ is represented by a cocycle \bar{u} in $F^*(A_3^*)$:

Let ρ be a cochain

$$\begin{split} &-\left[\tau_{2}\mid\varepsilon_{1}\mid\varepsilon_{1}^{9}\tau_{1}\right]-\left[\tau_{2}\mid\varepsilon_{1}^{10}\mid\tau_{1}\right]-\left[\varepsilon_{1}^{9}\tau_{2}\mid\varepsilon_{1}\mid\tau_{1}\right]\\ &+\left[\tau_{3}\mid\varepsilon_{1}\mid\tau_{1}\right]+\left[\varepsilon_{2}\mid\tau_{1}\mid\varepsilon_{1}^{9}\tau_{1}\right]+\left[\varepsilon_{2}\mid\varepsilon_{1}^{9}\tau_{1}\mid\tau_{1}\right]\\ &+\left[\varepsilon_{1}^{9}\varepsilon_{2}\mid\tau_{1}\mid\tau_{1}\right]-\left[\varepsilon_{3}\mid\tau_{1}\mid\tau_{1}\right]-\left[\tau_{3}\mid\varepsilon_{1}^{2}\mid\tau_{0}\right]\\ &-\left[\varepsilon_{3}\mid\varepsilon_{1}\tau_{0}\mid\tau_{1}\right]+\left[\varepsilon_{3}\mid\varepsilon_{1}\mid\tau_{1}\tau_{0}\right]-\left[\varepsilon_{3}\mid\varepsilon_{1}\tau_{1}\mid\tau_{0}\right]\\ &+\left[\varepsilon_{3}\mid\varepsilon_{1}^{2}\tau_{0}\mid\tau_{0}\right]-\left[\varepsilon_{3}^{2}\mid\varepsilon_{1}\tau_{1}\mid\tau_{1}\right]+\left[\varepsilon_{3}^{2}\mid\varepsilon_{1}^{2}\tau_{1}\mid\tau_{0}\right]\\ &-\left[\varepsilon_{3}^{2}\mid\varepsilon_{1}^{2}\mid\tau_{1}\sigma_{0}\right]+\left[\varepsilon_{1}^{9}\tau_{2}\mid\varepsilon_{1}^{2}\mid\tau_{0}\right]+\left[\tau_{2}\mid\varepsilon_{1}^{11}\mid\tau_{0}\right]\\ &+\left[\varepsilon_{1}^{2}\mid\varepsilon_{1}^{2}\mid\varepsilon_{1}^{9}\tau_{0}\right]-\left[\varepsilon_{1}^{9}\varepsilon_{2}\mid\varepsilon_{1}\mid\tau_{1}\tau_{0}\right]+\left[\varepsilon_{1}^{9}\varepsilon_{2}\mid\varepsilon_{1}^{10}\tau_{0}\mid\tau_{1}\right]\\ &+\left[\varepsilon_{1}^{9}\varepsilon_{2}\mid\varepsilon_{1}^{10}\tau_{1}\mid\tau_{0}\right]-\left[\varepsilon_{2}\mid\varepsilon_{1}^{10}\mid\tau_{1}\tau_{0}\right]+\left[\varepsilon_{3}\mid\varepsilon_{1}^{10}\tau_{0}\mid\tau_{1}\right]\\ &+\left[\varepsilon_{2}\mid\varepsilon_{1}^{10}\tau_{1}\mid\tau_{0}\right]-\left[\varepsilon_{2}\mid\varepsilon_{1}^{1}\mid\varepsilon_{1}^{9}\tau_{1}\tau_{0}\right]+\left[\varepsilon_{3}\mid\varepsilon_{1}^{10}\tau_{0}\mid\varepsilon_{1}^{9}\tau_{1}\right]\\ &+\left[\varepsilon_{2}\mid\varepsilon_{1}^{10}\tau_{1}\mid\varepsilon_{0}\right]-\left[\varepsilon_{2}\mid\varepsilon_{1}^{1}\mid\varepsilon_{1}^{9}\tau_{1}\tau_{0}\right]+\left[\varepsilon_{3}\mid\varepsilon_{1}^{10}\tau_{0}\mid\varepsilon_{1}^{9}\tau_{1}\right]\\ &+\left[\varepsilon_{2}\mid\varepsilon_{1}^{10}\tau_{1}\mid\varepsilon_{0}\right]-\left[\varepsilon_{2}\mid\varepsilon_{1}\mid\varepsilon_{1}^{1}\tau_{1}\tau_{0}\mid\varepsilon_{1}^{9}\right]-\left[\varepsilon_{1}^{9}\varepsilon_{2}\mid\varepsilon_{1}^{2}\tau_{0}\mid\tau_{0}\right] \end{split}$$

$$\begin{split} &-\left[\left.\boldsymbol{\varepsilon}_{2}\right|\,\boldsymbol{\varepsilon}_{1}^{11}\boldsymbol{\tau}_{0}\,|\,\boldsymbol{\tau}_{0}\right]-\left[\left.\boldsymbol{\varepsilon}_{2}\right|\,\boldsymbol{\varepsilon}_{1}^{2}\boldsymbol{\tau}_{0}\right|\,\boldsymbol{\varepsilon}_{1}^{9}\boldsymbol{\tau}_{0}\right]+\left[\left.\boldsymbol{\varepsilon}_{1}^{12}\right|\,\boldsymbol{\varepsilon}_{1}\boldsymbol{\tau}_{1}\,|\,\boldsymbol{\tau}_{1}\right]\\ &+\left[\left.\boldsymbol{\varepsilon}_{1}^{3}\right|\,\boldsymbol{\varepsilon}_{1}^{10}\boldsymbol{\tau}_{1}\,|\,\boldsymbol{\tau}_{1}\right]+\left[\left.\boldsymbol{\varepsilon}_{1}^{3}\right|\,\boldsymbol{\varepsilon}_{1}\boldsymbol{\tau}_{1}\,|\,\boldsymbol{\varepsilon}_{1}^{9}\boldsymbol{\tau}_{1}\right]-\left[\left.\boldsymbol{\varepsilon}_{1}^{12}\right|\,\boldsymbol{\varepsilon}_{1}^{2}\boldsymbol{\tau}_{1}\,|\,\boldsymbol{\tau}_{0}\right]\\ &+\left[\left.\boldsymbol{\varepsilon}_{1}^{12}\right|\,\boldsymbol{\varepsilon}_{1}^{2}\,|\,\boldsymbol{\tau}_{1}\boldsymbol{\tau}_{0}\right]-\left[\left.\boldsymbol{\varepsilon}_{1}^{3}\right|\,\boldsymbol{\varepsilon}_{1}^{11}\boldsymbol{\tau}_{1}\,|\,\boldsymbol{\tau}_{0}\right]+\left[\left.\boldsymbol{\varepsilon}_{1}^{3}\right|\,\boldsymbol{\varepsilon}_{1}^{11}\,|\,\boldsymbol{\tau}_{1}\boldsymbol{\tau}_{0}\right]\\ &+\left[\left.\boldsymbol{\varepsilon}_{1}^{3}\right|\,\boldsymbol{\varepsilon}_{1}^{2}\,|\,\boldsymbol{\varepsilon}_{1}^{9}\boldsymbol{\tau}_{1}\boldsymbol{\tau}_{0}\right]-\left[\left.\boldsymbol{\varepsilon}_{1}^{3}\right|\,\boldsymbol{\varepsilon}_{1}^{2}\boldsymbol{\tau}_{1}\,|\,\boldsymbol{\varepsilon}_{1}^{9}\boldsymbol{\tau}_{0}\right]+\left[\left.\boldsymbol{\varepsilon}_{1}^{15}\right|\,\boldsymbol{\tau}_{0}\,|\,\boldsymbol{\tau}_{0}\right]\\ &+\left[\left.\boldsymbol{\varepsilon}_{1}^{6}\right|\,\boldsymbol{\varepsilon}_{1}^{9}\boldsymbol{\tau}_{0}\,|\,\boldsymbol{\tau}_{0}\right]+\left[\left.\boldsymbol{\varepsilon}_{1}^{6}\right|\,\boldsymbol{\tau}_{0}\,|\,\boldsymbol{\varepsilon}_{0}^{9}\boldsymbol{\tau}_{0}\right].\end{split}$$

By tedious but routine calculations, we have

.
$$\delta$$
 $(\rho) = \bar{u}$. $[\xi_1^9] + ([\xi_2^3 \mid \xi_1^3] - [\xi_1^9 \mid \xi_1^6])$. $[\tau_0 \mid \tau_0]$.

Then we have the relation a).

Last we consider the relation e). Let ρ be the cochain defined in proof of a) and

$$\begin{array}{l} \mu \ \ \text{be} \ \ \left[\left. \tau_{\,0} \, \right| \, \xi_{\,1}^{\,2} \, \tau_{\,0} \, \right| \, \xi_{\,1} \right] \, - \, \left[\left. \tau_{\,0} \, \right| \, \xi_{\,1}^{\,2} \, \right| \, \xi_{\,1} \, \tau_{\,0} \right] \, + \, \left[\left. \tau_{\,0} \, \right| \, \xi_{\,1} \, \tau_{\,0} \, \right| \, \xi_{\,1}^{\,2} \right] \\ - \, \left[\left. \tau_{\,0} \, \right| \, \xi_{\,1} \, \right| \, \xi_{\,1}^{\,2} \, \tau_{\,0} \right] \, + \, \left[\left. \tau_{\,0} \, \right| \, \xi_{\,1}^{\,2} \, \right| \, \tau_{\,1} \right] \, - \, \left[\left. \xi_{\,1} \, \tau_{\,0} \, \right| \, \xi_{\,1} \, \right| \, \tau_{\,1} \right] \\ + \, \left[\left. \xi_{\,1} \, \tau_{\,0} \, \right| \, \xi_{\,1}^{\,2} \, \right| \, \tau_{\,0} \right] \, + \, \left[\left. \tau_{\,1} \, \right| \, \xi_{\,1} \, \right| \, \tau_{\,1} \right] \, - \, \left[\left. \tau_{\,1} \, \right| \, \xi_{\,1}^{\,2} \, \right| \, \tau_{\,0} \right] \, \cdot \\ \end{array}$$

Then, by the tedious but routine calculations, we see that $ub_{02}h_0 \in H^{6,78}(A_3)$ is represented by a cocycle $\overline{ub_{02}h_0}$ in $F^*(A_3^*)$:

$$\overline{ub_{02}h_0} = u \cdot \overline{b_{02}} \cdot [\xi_1] + u \cdot \overline{b_{11}} \cdot [\xi_2] - \rho \cdot \overline{b_{01}} \cdot [\xi_1]
+ ([\xi_2^3] \xi_1^3] - [\xi_1^9] \xi_1^6]) \cdot \mu \cdot [\xi_1] \cdot$$

Let v and G be

$$\begin{split} &-\left[\begin{smallmatrix}\tau_{3} \mid \, \xi_{1} \mid \, \tau_{1}\end{smallmatrix}\right] + \left[\begin{smallmatrix}\xi_{3} \mid \, \tau_{1} \mid \, \tau_{1}\end{smallmatrix}\right] + \left[\begin{smallmatrix}\tau_{3} \mid \, \xi_{1}^{2} \mid \, \tau_{0}\end{smallmatrix}\right] \\ &+\left[\begin{smallmatrix}\xi_{3} \mid \, \xi_{1}\tau_{0} \mid \, \tau_{1}\end{smallmatrix}\right] - \left[\begin{smallmatrix}\xi_{3} \mid \, \xi_{1} \mid \, \tau_{1}\tau_{0}\end{smallmatrix}\right] + \left[\begin{smallmatrix}\xi_{3} \mid \, \xi_{1}\tau_{1} \mid \, \tau_{0}\end{smallmatrix}\right] \\ &-\left[\begin{smallmatrix}\xi_{3} \mid \, \xi_{1}^{2}\tau_{0} \mid \, \tau_{0}\end{smallmatrix}\right] + \left[\begin{smallmatrix}\xi_{3}^{3} \mid \, \xi_{1}\tau_{1} \mid \, \tau_{1}\end{smallmatrix}\right] - \left[\begin{smallmatrix}\xi_{3}^{3} \mid \, \xi_{1}^{2}\tau_{1} \mid \, \tau_{0}\end{smallmatrix}\right] \\ &+\left[\begin{smallmatrix}\xi_{3}^{3} \mid \, \xi_{1}^{2} \mid \, \tau_{1}\tau_{0}\end{smallmatrix}\right] + \left[\begin{smallmatrix}\xi_{1}^{9} \mid \, \xi_{2} \mid \, \xi_{1}\tau_{1}\tau_{0}\end{smallmatrix}\right] \end{split}$$

and

 $-\left[\xi_{1}^{9}\right] \cdot \rho - \nu \cdot \left[\xi_{1}^{9}\right] , \text{ respectively. Let } \overline{c} \text{ be the representative } \left[\xi_{2}^{3} \mid \xi_{1}^{3}\right] - \left[\xi_{1}^{9} \mid \xi_{1}^{6}\right] \text{ of } c \in H^{2,60}(A_{3}). \text{ Since } \delta \in \left[\xi_{1}^{9}\right] \cdot \rho + \nu \cdot \left[\xi_{1}^{9}\right] \right] = -\left[\xi_{1}^{9}\right] \cdot \overline{c} \cdot \left[\tau_{0} \mid \tau_{0}\right] - \overline{c} \cdot \left[\tau_{0} \mid \tau_{0}\right] \mid \xi_{1}^{9}\right] \text{ in } F^{*}(A_{3}^{*}) \text{ and } \delta_{3}(G) = -a_{0}^{2}h_{2}c \text{ in the May spectral sequence, the cochain } \overline{G} \text{ is a representative of } G \text{ in some sense. It is easy to see that}$

 $\equiv [\xi_1^9] \cdot \overline{ub_{02}h_0} - G \cdot \overline{b_{01}} \cdot [\xi_1]$ modulo terms which have the May's weight (the weight associated to May spectral sequence) less than 5 or have

the May's weight 5 and the weight, associated to the spectral sequence defined in (9), greater than 1. For the dimensional and filtrational reasons, we have the relation

 $h_2 \cdot ub_{02}h_0 = Gb_{01}h_0 + \varepsilon \ a_0^6h_3$, where $\varepsilon \in Z_3$. Since $d_2(h_2) = a_0b_{11}$ (Theorem 2.1.) and $d_2(h_0w) = h_0ub_{02}$ (Theorem 2.4.) in the mod 3 Adams spectral sequence, we have $d_2(h_2h_0w) = h_2 \cdot h_0ub_{02} = -h_0b_{01}G - \varepsilon \cdot a_0^6h_3$, up to sign. Then we have a non-zero differential $d_2(h_2w)$. By the dimensional considerations, we have $d_2(h_2w) = b_{01}G$, up to sign. Then we have $\varepsilon = 0$ and therefore we have the relation f).

§ 2. Some known results on the mod 3 Adms spectral sequence

From now on we neglect the non-zero coefficient of the differentials and the relations in the mod 3 Adams spectral sequence.

The following five theorems for the differentials and elements surviving to E_{∞} were verified by using the results of the stable homotopy groups of sphere and the statement that $E_{\rm r}$ is a differential algebra by J. P. May [5] (in the range $t-s \leq 32$ and partially in the range $t-s \geq 33$) and by S. Oka the rest.

THEOREM 2. 1. (H. H. Gershenson [2]. A. Liulevicius [4]. R.J. Milgram [7]. N. Shimada and T. Yamanoshita [11]. H. Toda [12])

$$d_2(h_i) = a_0 b_{i-1}$$
 for $i \ge 1$,

THEOREM 2. 2. (H.Toda [12,13,14,15])

a).
$$d_2(g_2) = g_1 b_{01}$$
 ([12])

b).
$$d_2(u) = a_1 b_{01}$$
 ([12])

c).
$$d_5(b_{11}) = h_0 b_{01}^3$$
 ([14])

d).
$$d_5(h_1b_{02}) = b_{01}^3k$$
 ([15])

e).
$$d_6(e_1) = b_{01}^6$$
 ([13]).

THEOREM 2. 3. (J.P.May [5])

a).
$$d_3(a_0^4h_2) = b_{01}a_1^2$$

b).
$$d_2(a_1^j u) = a_1^{j+1} b_{01}$$
 for $0 \le j$, $j \ne 1 \mod 3$.

THEOREM 2. 4. (S.Oka [10])

a).
$$d_2(c) = a_0 h_1 b_{02}$$

b),
$$d_3(a_0^2c) = h_0b_{01}^2b_{11}$$

c),
$$d_2(g_2a_2) = e_2$$

d),
$$d_2(f) = h_0 b_{01}^3 b_{11}$$

e).
$$d_3(a_0^8a_2u) = b_{01}a_1^5$$

f).
$$d_2(a_2u) = \ell$$

g).
$$d_2(h_0w) = h_0ub_{02}$$

h),
$$d_2(g_1w) = g_1ub_{02}$$

Remark. All other non-zero differentials in the range $t-s \le 77$ are easily determined by the differentials listed in the above four theorems and the statement that E_r is a differential algebra.

Theorem 2, 5, (J. P. May [5] and S. Oka) E_{∞} for $t-s \leq 76$ (and corresponding generators of Π_{*}^{s} (So:3)) are given in Table A.

Table A.

t-s	πs _* (S0:3)	survivor (corresponding generator)
0	Z	a_0^i (ι)
1	0	•
2	0	
3	Z_3	$h_0 (\alpha_1)$
4	0	•
5	0	
6	0	
7	$oldsymbol{Z_3}{0}$	$g_1(\alpha_2)$
8	0	
9	0	
10	Z_3	$b (\beta_1)$
11	Z_9	$a_0^{}h_1^{}$, $a_0^2h_1^{}$ (α_3^{\prime})
12	0	
13	Z_3	$h_0b (\alpha_1\beta_1)$
14	0	
15	Z_3	$h_{0}a_{1}(\alpha_{4})$
16	0	• • •
17	0	
18	0	
19	Z_3	$g_1 a_1 (\alpha_5)$
20	Z_3	$b^2 \left(\beta_1^2\right)$
21	0	
22	0	
23	$Z_3 + Z_9$	$h_0b^2 (\alpha_1\beta_1^2), a_0^2u, a_0^3u (\alpha_6')$
24	0	V - 1-1 V - U - U-
25	0	

Table A. (continued)

t-s	π ^s _* (S ⁰ :3)	survivor (corresponding generator)
26	Z_3	$k(\beta_2)$
27	Z_3	$h_0 a_1^2 \left(\alpha_{\gamma}\right)$
28	0	
29	Z_3	$h_0k (\alpha_1\beta_2)$
30	Z_3	b ³ (β ³)
31	Z_3	$g_1 a_1^2 \left(\alpha_8\right)$
32	0	
33	, Õ	
34	0	,
35	Z_{27}	$a_0^6h_2^{}$, $a_0^7h_2^{}$, $a_0^8h_2^{}$ ($\alpha_9^{'}$)
36	$Z_3^{\prime\prime}$	$bk (\beta_1 \beta_2)$
37	Z_3	$h_0 b_{11} (\epsilon')$
38	Z_3	$h_0 h_2 (\epsilon_1)$
39	$Z_3 + Z_3$	$h_0 bk \ (\alpha_1 \beta_1 \beta_2), \ h_0 a_1^3 \ (\alpha_{10})$
40	Z_3	b ⁴ (β ⁴ ₁)
41	0	
42	Z_3	$g_1 h_2 (\epsilon_2)$
43	Z_3	$g_1a_1^3$ (α_{11})
44	0	,
45	Z_9	bh_2 , a_0bh_2 (φ)
46	Z_3	$b^2k \left(\beta_1^2\beta_2\right)$
47		h_0bb_{11} $(\beta_1\epsilon')$, $a_0^2a_1^2u$, $a_0^3a_1^2u$ (α'_{12})
48	0	
49	Z_3	$h_0b^2k \left(\alpha_1\beta_1^2\beta_2\right)$
50	Z_3	$b^5 \left(\beta \frac{5}{1}\right)$
51 52	Z_3	$h_0a_1^4 \left(\alpha_{13}\right)$
	Z_{3}	$k^2 \left(\beta \frac{2}{2}\right)$
53	0	
54	0	
55	$Z_{3} + Z_{3}$	h_0k^2 $(\alpha_1\beta_2^2)$, $g_1a_1^4$ (α_{14})
56	0	
57	0	
58	0	

t-s	π _* (S ⁰ :3)	survivor (corresponding generator)
59	Z_{g}	$a_0^2 a_1^3 u$, $a_0^3 a_1^3 u$ (α'_{15})
60	0	
61	0	
62	Z_3	$bk^2 \left(\beta_1\beta_2^2\right)$
63	$Z_3 \ Z_3$	$h_0 a_1^5 \left(\alpha_{16}^{} \right)$
64	0	
65	$Z_3^{}$	$h_0 b k^2 (\alpha_1 \beta_1 \beta_2^2)$
66	0	
67	Z_3	$g_1a_1^5$ (α_{17})
68	Z_3	$b_{11}^2(\lambda)$
69	0	
70	0	
71	$Z_{\overline{27}}$	$a_0^{10}a_2u$, $a_0^{11}a_2u$, $a_0^{12}a_2u$ ($\alpha_{18}^{"}$)
72	$Z_3^{}$	$b^2k^2 \ (\beta \ _1^2 \beta \ _2^2)$
73	0	
74	$Z^{}_3$	$ka_2(\beta_5)$
75	$Z_3 + Z_3 + Z_3$	h_0bk_{02} (μ), $h_0b^2k^2$ ($\alpha_1\beta_1^2\beta_2^2$), $h_0a_1^6$ (α_{19})
76	0	

Table A' (continued)

The matric Massey products in algebraic spectral sequences were studied by J.P.May (6). We quote some of his results which we use in this paper.

Theorem 2. 6. (J.P.May). Let $\langle v^1, v^2, v^3 \rangle$ be defined in E_{r+1} term of the May spectral sequence. Assume that $v^i \in E_{r+1}^{p_i, q_i, t_i}$ and that v^i converges to w^i , where $\langle w^1, w^2, w^3 \rangle$ is defined in H^{**} (A_3). Assume further that the following condition (*) is satisfied.

(*) If
$$(p,q,t) = (p_{i+1} + p_i, q_{i+1} + q_i, t_{i+1} + t_i)$$
, $i=1$ or 2, then $E_{r+u+1}^{p-r-u} q+r+u-1, t$ ($E_{r+u+1} \infty$ for $u \ge 0$.

Then any element of $\langle v^1, v^2, v^3 \rangle$ is a permanent cocycle which converges to an element of $\langle w^1, w^2, w^3 \rangle$.

The matric Massey products in the Adams spectral sequence were studied by R.M.F.Moss [8] and A.F.Lawrence [3]. We quote some of their results which we use in this paper.

THEOREM 2. 7, (R.M.F.Moss).